

Linear Systems, Black Boxes, and Beyond

Homework #1 (2010-2011), Answers

Q1: Fourier transforms, derivatives, and integrals

Setup is $\hat{s}(\omega) = \int_{-\infty}^{\infty} s(t)e^{-i\omega t} dt$, with $s(t) = \frac{1}{2\pi} \int_{-\infty}^{\infty} \hat{s}(\omega)e^{+i\omega t} d\omega$.

A. For $q(t) = \frac{d}{dt}s(t)$, find $\hat{q}(\omega)$.

Using the “synthesis” integral,

$$q(t) = \frac{d}{dt}s(t) = \frac{d}{dt} \left(\frac{1}{2\pi} \int_{-\infty}^{\infty} \hat{s}(\omega)e^{+i\omega t} d\omega \right) = \frac{1}{2\pi} \int_{-\infty}^{\infty} \hat{s}(\omega) \frac{d}{dt} (e^{+i\omega t}) d\omega = \frac{1}{2\pi} \int_{-\infty}^{\infty} \hat{s}(\omega) (i\omega e^{+i\omega t}) d\omega.$$

So the coefficient of $e^{i\omega t}$ in $q(t) = \frac{d}{dt}s(t)$ is $\hat{q}(\omega) = i\omega\hat{s}(\omega)$.

B. For $q_n(t) = \frac{d^n}{dt^n}s(t)$, find $\hat{q}_n(\omega)$.

Iterating part A: $\hat{q}_n(\omega) = i\omega\hat{q}_{n-1}(\omega)$, so $\hat{q}_n(\omega) = (i\omega)^n \hat{s}(\omega)$.

C. For $z(t) = \int_{-\infty}^t s(\tau)d\tau$, find $\hat{z}(\omega)$.

Since $s(t) = \frac{dz}{dt}$, we can use part A: $\hat{s}(\omega) = i\omega\hat{z}(\omega)$, so, except possibly at $\omega = 0$, $\hat{z}(\omega) = \frac{\hat{s}(\omega)}{i\omega}$

D. Apply C to $s(t) = \delta(t)$ to find a function whose Fourier transform, except possibly at 0, is

$$\frac{1}{i\omega}.$$

Since the Fourier transform of the delta-function is 1 everywhere, the integral of the delta-function, $h(t) = \int_{-\infty}^t \delta(\tau)d\tau$ has the required Fourier transform $\frac{1}{i\omega}$. Since the delta-function is an

infinitesimally narrow peak with a unit area, the integral evaluates as $h(t) = \begin{cases} 1, t > 0 \\ 0, t < 0 \end{cases}$. This is the

“Heaviside step function.” Its value at zero, which is formally undefined, is irrelevant for most purposes.

Q2: Fourier transforms and moments

Setup is $\hat{s}(\omega) = \int_{-\infty}^{\infty} s(t)e^{-i\omega t} dt$, with $s(t) = \frac{1}{2\pi} \int_{-\infty}^{\infty} \hat{s}(\omega)e^{+i\omega t} d\omega$, but now we are thinking of s as a probability distribution.

A. Write the normalization condition $\int_{-\infty}^{\infty} s(t)dt = 1$ in terms of $\hat{s}(\omega)$.

Since $e^{i\omega t} = 1$ for $\omega = 0$, $\hat{s}(0) = \int_{-\infty}^{\infty} s(t)dt$, so the normalization condition is $\hat{s}(0) = 1$.

B. Write the mean (first moment) $\langle t \rangle = \int_{-\infty}^{\infty} ts(t)dt$ in terms of $s'(\omega) = \frac{d}{d\omega} \hat{s}(\omega)$.

Since $\hat{s}'(\omega) = \int_{-\infty}^{\infty} s(t) \frac{d}{d\omega} e^{-i\omega t} dt = \int_{-\infty}^{\infty} s(t)(-it)e^{-i\omega t} dt$, it follows that $\hat{s}'(0) = \int_{-\infty}^{\infty} s(t)(-it)dt$ and

that $\int_{-\infty}^{\infty} ts(t)dt = i\hat{s}'(0)$.

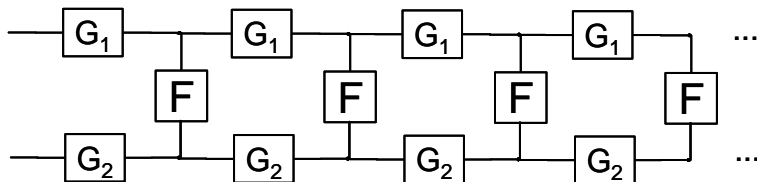
C. Write the variance (second moment) $\langle (t - \langle t \rangle)^2 \rangle = \langle t^2 \rangle - \langle t \rangle^2 = \int_{-\infty}^{\infty} t^2 s(t)dt - \left(\int_{-\infty}^{\infty} ts(t)dt \right)^2$ in

terms of $s'(\omega) = \frac{d}{d\omega} \hat{s}(\omega)$ and $s''(\omega) = \frac{d^2}{d\omega^2} \hat{s}(\omega)$.

As in part B, $\hat{s}''(\omega) = \int_{-\infty}^{\infty} s(t) \frac{d^2}{d\omega^2} e^{-i\omega t} dt = \int_{-\infty}^{\infty} s(t)(-t^2)e^{-i\omega t} dt$, so $\int_{-\infty}^{\infty} t^2 s(t)dt = -\hat{s}''(0)$.

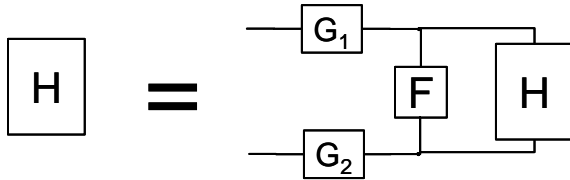
So $\int_{-\infty}^{\infty} t^2 s(t)dt - \left(\int_{-\infty}^{\infty} ts(t)dt \right)^2 = -\hat{s}''(0) - (i\hat{s}'(0))^2 = -\hat{s}''(0) + (\hat{s}'(0))^2$.

Q3: The half-infinite cable (repeating indefinitely to the right)



This is to be viewed as a network of resistors and capacitors. Calculate the impedance of the system (input applied across terminals at left) in terms of the impedances $F(\omega)$, $G_1(\omega)$, and $G_2(\omega)$ for F , G_1 , and G_2 .

Hint: Let the composite system be H . Note the following, and then write an equation for $H(\omega)$.



The impedance of the composite system on the left is a series combination of three components: G_1 , the parallel combination of F and H , and G_2 . Therefore its impedance is

$G_1(\omega) + \frac{F(\omega)H(\omega)}{F(\omega) + H(\omega)} + G_2(\omega)$. Since (as the hint indicates) this is equivalent to the entire half-

infinite cable, $H(\omega) = G_1(\omega) + \frac{F(\omega)H(\omega)}{F(\omega) + H(\omega)} + G_2(\omega)$. Solving for $H(\omega)$ yields

$H(\omega)^2 - G(\omega)H(\omega) - G(\omega)F(\omega) = 0$, where $G(\omega) = G_1(\omega) + G_2(\omega)$, or,

$$H(\omega) = \frac{G(\omega) + \sqrt{G(\omega)^2 + 4F(\omega)G(\omega)}}{2}.$$

Note concerning the continuum limit: This corresponds to allowing each subunit to represent progressively less and less length. Then F has units of impedance/cm (and increases as the segment shortens), and G has units of impedance-cm (and decreases as the segment shortens). In this limit, $H(\omega) \approx \sqrt{F(\omega)G(\omega)}$. This enables one to calculate the ‘‘cable length’’ λ , which is the distance required for the transmembrane current to fall by a factor of e . To do this, note that total

transmembrane current I_{total} is $\int_0^{\infty} e^{-x/\lambda} dx = \lambda$ times the current per unit length I_{peak} at the

injection site, but also, I_{total} / I_{peak} is inversely proportional to the total cable impedance $H(\omega)$,

divided by the impedance per unit length, $F(\omega)$. So $\lambda = \frac{H(\omega)}{F(\omega)} = \sqrt{\frac{G(\omega)}{F(\omega)}}$.

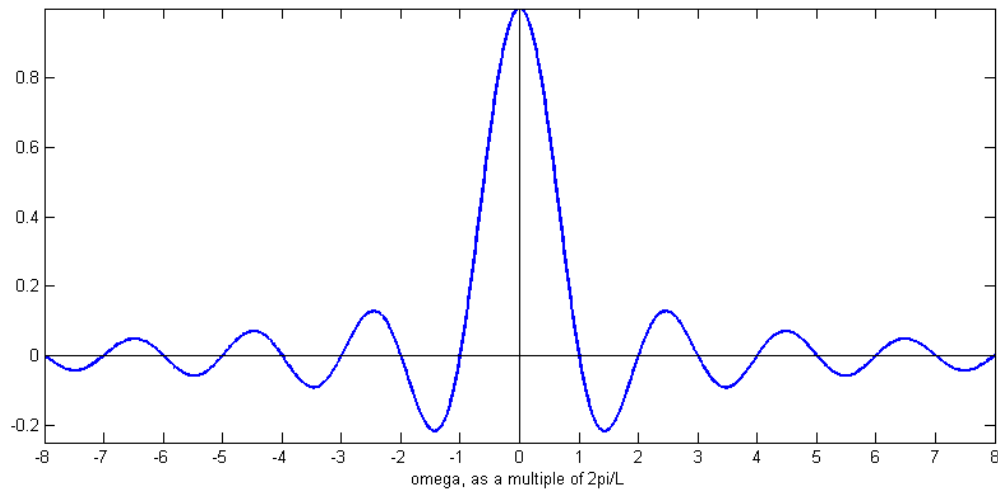
Q4. Boxcar smoothing

Boxcar smoothing refers to convolution with the function $s(t)$, where $s(t) = \begin{cases} \frac{1}{L}, & |t| \leq L/2 \\ 0, & |t| > L/2 \end{cases}$. Find

its Fourier transform. What does it look like? Is this a good way to smooth?

$$\hat{s}(\omega) = \int_{-\infty}^{\infty} s(t)e^{-i\omega t} dt = \frac{1}{L} \int_{-L/2}^{L/2} e^{-i\omega t} dt = \frac{1}{-i\omega L} e^{-i\omega t} \Big|_{-L/2}^{L/2} = \frac{e^{i\omega L/2} - e^{-i\omega L/2}}{i\omega L} = \frac{\sin(\omega L/2)}{(\omega L/2)}.$$

This (the “sinc” function) has a peak of 1 at $\omega = 0$, and descends in an envelope proportional to $1/|\omega|$ away from zero. There are zeros at $\omega = 2\pi k/L$, for $k \neq 0$. The center lobe (at $\omega = 0$) is positive, but the adjacent lobes ($\frac{2\pi}{L} < |\omega| < \frac{4\pi}{L}$) are negative. So one problem with using this as a smoothing function is that it inverts the phase of non-negligible frequency components.



```
>> x=[-8:0.01:8];
>> y=sinc(pi*x);
>> plot(x,y)
>> hold on;
>> plot([-8 8],[0 0],'k')
>> plot([0 0],[-0.5 1],'k')
>> set(gca,'YLim',[-0.2 1])
>> set(gca,'YLim',[-0.25 1])
>> xlabel('omega, as a multiple of 2pi/L')
>> set(gca,'XTick',[-8:8])
```